

Research Article

Study of Triangular Fuzzy Hybrid Nanofluids on the Natural Convection Flow and Heat Transfer between Two Vertical Plates

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Received 3 September 2021; Revised 22 September 2021; Accepted 27 September 2021; Published 11 November 2021

Academic Editor: Ahmed Mostafa Khalil

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The prime objective of the current study is to examine the effects of third-grade hybrid nanofluid with natural convection utilizing the ferro-particle (Fe_3O_4) and titanium dioxide (TiO_2) and sodium alginate (SA) as a host fluid, flowing through vertical parallel plates, under the fuzzy atmosphere. The dimensionless highly nonlinear coupled ordinary differential equations are computed adopting the bvp4c numerical approach. This is an extremely effective technique with a low computational cost. For validation, it is found that as the volume fraction of ($\text{Fe}_3\text{O}_4 + \text{TiO}_2$) hybrid nanoparticles rises, so does the heat transfer rate. The current and existing results with their comparisons are shown in the form of the tables. The present findings are in good agreement with their previous numerical and analytical results in a crisp atmosphere. The nanoparticles volume fraction of Fe_3O_4 and TiO_2 is taken as uncertain parameters in terms of triangular fuzzy numbers (TFNs) [0, 0.05, 0.1]. The TFNs are controlled by α -cut and the variability of the uncertainty is studied through triangular membership function (MF).

1. Introduction

Researchers have been attracted by natural convection (NC) flow because of its numerous uses in engineering and scientific problems like heat exchangers, building ventilation, insulation, solar energy collection, refrigeration, nuclear waste repositories, petroleum reservoirs geothermal systems, and chemical catalytic reactors. Convection is used significantly in the manufacturing of solar panels, microstructures during the cooling of molten metals, and free air cooling without the need for fans in real-world applications. Various researchers have looked into the NC-based flow of non-Newtonian and Newtonian fluids between two infinite parallel vertical plates such as Bruce and Na [1] who investigated the heat transfer of NC between vertical flat plates

using non-Newtonian Powell–Eyring fluids. Later on, Rajagopal and Na [2] studied the extensive thermodynamic analysis on fundamental functions. The influences of the third-grade non-Newtonian fluid on heat transfer (HT) were examined by Ziabakhsh and Domairry [3] through the homotopy analysis method (HAM). Using the least square method (LSM), Maghsoudi et al. [4] inspected the NC flow of third-grade fluid between two infinite vertical flat plates with a porous media. Mansoor et al. [5] studied the natural convective flow between two vertical plates with the help of the volume of parameter method (VPM) and Runge–Kutta method (RKM). They show that VPM is better than RKM. Some researchers have explored different flows of fluids between vertical parallel plates analytically and numerically [6, 7].

Because of the escalation in energy prices, HT management is extremely important in energy systems. So the nanofluids are the mixtures of liquid and nanoparticles which are used to improve the rate of heat transfer. The performance of nanoparticles in a heat transfer mechanism is excellent as compared to regular fluid. This is due to the dispersed ultrafine particles boosting the thermal conductivity of the fluid and therefore increasing their energy transfer competency. In this way, convective HT of nanofluid is a challenging problem. When a non-Newtonian fluid is moving in a structure, natural convection analysis is a difficult task. The nanoparticle adds to the base fluid and then heat transfer may be increased and this added nanoparticle in the base fluid is named nanofluid [8, 9]. Numerous literatures [10, 11] disclose the low volume fractions (1–4 volume %) for better performance of thermal conductivity of the fluids. We can utilise nanoparticle concentrations of greater than 20% [12]. When two or more distinct nanoparticles are added to the base fluid, it is referred to as “hybrid nanofluid,” also the thermal conductivity is greater than as compared to nanofluid and regular fluid. In the field of heat transfer, hybrid nanofluids have received a lot of attention such as nuclear system cooling, drug reduction, automobile radiators, thermal storage, welding, electronic cooling, solar heating, lubrication, the coolant in machining, generator cooling, defence, biomedical, heating, and refrigeration, etc. Saqib et al. [13] studied the NC flow problem on Jaffry hybrid nanofluid using CNTs (single- and multiwall carbon nanotubes) with carboxy-methyl-cellulose (CMS) as a base fluid between two vertical parallel plates. Hatami and Ganji [14] applied the differential transform method (DTM) to investigate the NC flow of sodium alginate (SA) as a host fluid and silver (Ag) and copper (Cu) as nanofluids between two vertical parallel plates. Maghsoudi et al. [15] investigated natural convective, thermal radiation, HT, and magnetic field of the non-Newtonian nanofluid flow between two infinite vertical flat plates utilizing the Galerkin method (GM). Using the HAM, Rahmani et al. [16] explored the NC flow of non-Newtonian nanofluids between two vertical plates. They observed that HAM is better than the numerical RK method. The NC flow of non-Newtonian nanofluids between two vertical plates using the generalized decomposition method (GDM) was also studied by Kezzar et al. [17]. They observed that GDM is better than the numerical RK method. Biswal et al. [18, 19] used the HPM in an uncertain environment to examine the NC of nanofluid flow between two parallel plates. The volume fraction of nanoparticle was considered as TFN and also shows the fuzzy result is better than a crisp result. Gabli et al. [20] studied the NC flow of non-Newtonian ferro-particle (Fe_3O_4) nanofluids between two vertical plates with thermal radiation using the Adomian decomposition method (ADM). They observed that ADM is better than the RK-Feldberg-based shooting method. Devi and Devi [21] inspected the HT and flow problems of hydro-magnetic hybrid nanofluids ($\text{Al}_2\text{O}_3 + (\text{Cu}/\text{H}_2\text{O})$) through a stretched sheet.

Fluid flow with heat transfer is essential in science and engineering. Because of extensive physical properties such as chemical diffusion, magnetic effect, and heat transfer, governing fluid equations are converted into linear or nonlinear DEs. After controlling these physical issues, they are transformed into linear or nonlinear DEs. The solution of DEs is strongly affected by the physical problems with associated parameters and initial, geometry, coefficient, and boundary conditions. Then, these are not crisp due to the mechanical defect, experimental error, and measurement error, etc. In this scenario, fuzzy sets theory (FST) is a more accurate instrument than assuming genuine physical problems for getting a better understanding of the facts under investigation. To be more specific, FDEs are useful for decreasing uncertainty and determining the best way to define a physical problem with unknown parameters and initial and boundary conditions.

The FST was first presented by Zadeh [21] in 1965. FST is a fantastic approach for describing circumstances when information is unclear, imprecise, or uncertain. Later on, Dubois and Prade [22] developed arithmetic procedures on fuzzy numbers (FNs). The trapezoidal, triangular, and Gaussian FNs are three forms of FNs that may be classified. For thoroughness, we will look at TFNs now. The FN is a variable that has a range from 0 to 1. Each numerical value in the range is given a membership grade, with 0 being the lowest grade and 1 being the strongest possible grade. The information contained in crisp partial or ordinary differential equations models of dynamical systems is sometimes incomplete, imprecise, or ambiguous. FDEs are a useful approach for modelling dynamical systems with ambiguity or uncertainty. FNs or TFNs can be used to define this impreciseness or vagueness mathematically. Many researches have been conducted in recent years around the notion of FDEs. The fuzzy differentiability idea was established by Seikala [23] and Kaleva [24] and then they discussed fuzzy integration and differentiation. The FDEs were first presented in 1987 by Kandel and Byatt [25]. For the solution of FDEs, Buckley et al. [26] employed two methods: the extension principle and FNs. For continuous FDEs, Nieto [27] investigated the Cauchy problem. In [28], Lakshmikantham and Mohapatra investigated the initial value problems for FDEs. For the existence and uniqueness solution of FDE, Park and Hyo [29] employed the successive approximation approach. The geometric approach for solving a system of FDEs was devised by Gasilov et al. [30]. The system of FDEs with TFNs was investigated by Nizami et al. [31]. Salahsour et al. [32] used FDE and TFNs to investigate the fuzzy alley effect and the fuzzy logistic equation.

In addition, numerous scholars have used FST to achieve well-known findings in commerce and science, for example, in bank account model [33], population dynamics model [34, 35], bacteria culture model [36], HIV model [37], growth model [38], computational biology [39], modelling hydraulic [40], predator-prey model [41], quantum optics and gravity [42], decay model [43], model of friction [44], civil engineering [45], Laplace transform [46], integro-

differential equation [47], dengue virus model [48], chemostat model [49], and giving up smoking model [50]. The adaptive fuzzy controller for uncertain fractional-order nonlinear systems was proposed by Liu et al. [51], and it was used to evaluate an unknown nonlinear function. Very recently, Nadeem et al. [52] recently investigated MHD and ohmic heating on the third-grade fluid in an inclined channel under the fuzzy environment. The triangle membership function was used to discuss the uncertainty.

Inspired by the earlier investigations, the goal of this paper is to use the numerical scheme bvp4c to analyze the hybrid nanofluid flow between two vertical parallel plates in a fuzzy environment. The sodium alginate (SA) is the host fluid, while the hybrid nanoparticles are Fe_3O_4 and TiO_2 . The impacts of the Eckert number, Prandtl number, and nanoparticle volume fraction on velocity and temperature profiles are studied. It has been detected that hybrid nanofluid enhanced the thermal efficiency of the base fluid rapidly as compared to the other fluid and nanofluids. Besides, after checking the accuracy of bvp4c so compare the results of existing work in the literature. The nanoparticle volume fraction has also been treated as an uncertain parameter in this investigation, with fuzzy numbers or triangular fuzzy numbers being used. The FDEs with α -cut approach are used to tackle the natural convection problem in a fuzzy environment.

The paper is arranged as follows. We discussed some fundamental preliminaries on FDEs in Section 2. The formation of the crisp problem is described in Section 3. The crisp problem is transformed into FDEs in Section 4. Section 5 contains an explanation of graphs and tables. Section 6 concludes with some closing remarks.

2. Preliminaries

Some fundamental definitions are given in this section.

Definition 1 (see [21, 52]). Fuzzy set is defined as a set of ordered pairs such that $\tilde{U} = \{(y, \mu_{\tilde{U}}(y)) : y \in X, \mu_{\tilde{U}}(y) \in [0, 1]\}$, where $\mu_{\tilde{U}}(y)$ is the membership function of \tilde{U} , X is the universal set, and mapping is defined as $\mu_{\tilde{U}}(y) : X \rightarrow [0, 1]$.

Definition 2 (see [21, 52]). α -cut or α -level of a fuzzy set \tilde{U} is a crisp set U_α and defined by $U_\alpha = \{(y/\mu_{\tilde{U}}(y)) \geq \alpha\}$, where $0 \leq \alpha \leq 1$.

Definition 3 (see [22, 52]). Let $\tilde{U} = (a_1, a_2, a_3)$ with membership function $\mu_{\tilde{U}}(y)$ which is called a TFN if

$$\mu_{\tilde{U}}(y) = \begin{cases} \frac{a_1 - y}{a_2 - a_1}, & \text{for } y \in [a_1, a_2], \\ \frac{y - a_3}{a_2 - a_3}, & \text{for } y \in [a_2, a_3], \\ 0, & \text{otherwise.} \end{cases} \quad (1)$$

The TFN with peak (or center) a_2 , left width $a_2 - a_1 > 0$, right width $a_3 - a_2 > 0$, and these TFNs are transformed into interval numbers through α -cut approach, is written as $\tilde{U} = [u_1(y; \alpha), u_2(y; \alpha)] = [a_1 + (-a_1 + a_2)\alpha, a_3 - (-a_2 + a_3)\alpha]$, where $\alpha \in [0, 1]$. The membership function is the building block of FST and it is defined by its membership function. TFNs $\tilde{U} = (a_1, a_2, a_3)$ and α -cut of membership function are shown in Figure 1. An arbitrary TFN satisfies the following conditions: (i) $u_1(y; \alpha)$ is an increasing function on $[0, 1]$. (ii) $u_2(y; \alpha)$ is a decreasing function on $[0, 1]$. (iii) $u_1(y; \alpha) \leq u_2(y; \alpha)$ on $[0, 1]$. (iv) $u_1(y; \alpha)$ and $u_2(y; \alpha)$ are bounded at $[0, 1]$, respectively. (v) If $u_1(y, \alpha) = u_2(y, \alpha) = u(y)$ where $u(y)$ is a crisp number.

Definition 4 (see [23, 25, 52]). Let I be a real interval. A mapping $\tilde{u}(y; \alpha) : I \rightarrow F$ is called a fuzzy process, defined as $\tilde{u}(y; \alpha) = [u_1(y; \alpha), u_2(y; \alpha)]$, $y \in I$ and $\alpha \in [0, 1]$. The derivative $(d\tilde{u}(y; \alpha)/dy) \in F$ of a fuzzy process $\tilde{u}(y; \alpha)$ is defined by $(d\tilde{u}(y; \alpha)/dy) = [(du_1(y; \alpha)/dy), (du_2(y; \alpha)/dy)]$.

Definition 5 (see [23, 25, 52]). Let $I \subseteq \mathbb{R}$, $\tilde{u}(y; \alpha)$ be a fuzzy valued function defined on I . Let $\tilde{u}(y; \alpha) = [u_1(y; \alpha), u_2(y; \alpha)]$ for all α -cut. Assume that $u_1(y; \alpha)$ and $u_2(y; \alpha)$ have continuous derivatives or are differentiable, for all $y \in I$ and α , then $(d\tilde{u}(y; \alpha)/dy) = [(du_1(y; \alpha)/dy), (du_2(y; \alpha)/dy)]_\alpha$. Similarly, we can define higher-order ordinary derivatives in the same way. An FN by an ordered pair of functions $[d\tilde{u}(y; \alpha)/dy]_\alpha$ and they satisfy the following conditions: (i) $(du_1(y; \alpha)/dy)$ and $(du_2(y; \alpha)/dy)$ are continuous on $[0, 1]$. (ii) $(du_1(y; \alpha)/dy)$ is an increasing function on $[0, 1]$. (iii) $(du_2(y; \alpha)/dy)$ is a decreasing function on $[0, 1]$. (iv) $(du_1(y; \alpha)/dy) \leq (du_2(y; \alpha)/dy)$ on $[0, 1]$.

3. Problem Formulation

In this proposed problem, Figure 2 portrays the main theme schematically. It consists of two vertical parallel flat plates separated by a distance $2h$ apart, in which there is a non-Newtonian fluid, which is flowing due to the free convection. The walls at $x = h$ and $x = -h$ are held at constant temperatures T_1 and T_2 , respectively, with $(T_1 > T_2)$. This difference of temperature causes the fluid near the walls at $x = -h$ to rise and the fluid near the wall $x = h$ to fall. The fluid is a non-Newtonian sodium alginate-based nanofluid containing Fe_3O_4 and TiO_2 hybrid nanoparticles. The base fluid and the hybrid nanoparticles are considered to be in thermal equilibrium, with no-slip between them. Some physical properties of the hybrid nanofluid are arranged in Table 1.

Using the above assumptions and Boussinesq approximation [14], the momentum and energy equations of the

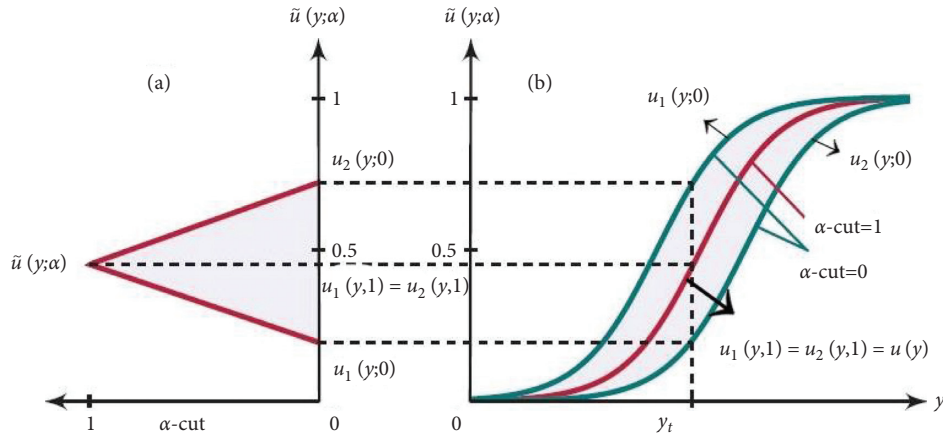


FIGURE 1: Membership functions of a TFN.

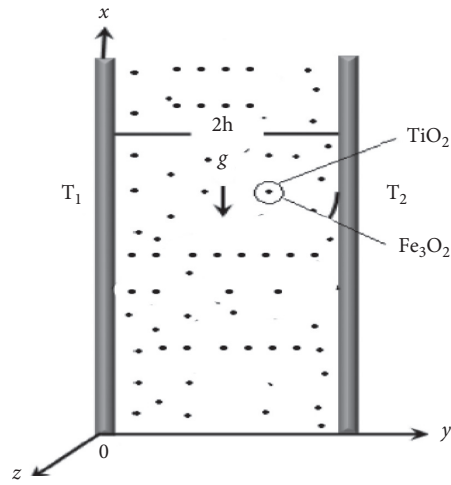


FIGURE 2: Flow geometry.

TABLE 1: Thermo-physical properties of base fluids and hybrid nanoparticles [14, 20].

Materials	ρ (kg/m ³)	C_p (J/kg ⁻¹ k ⁻¹)	K (W/m)	β_T (k ⁻¹)
Sodium alginate (SA)	989.0	4175.0	0.6376.0	99.0
Ferro-particle (Fe ₃ O ₄)	5180.0	670.0	9.70	1.18×10^{-5}
Titanium dioxide (TiO ₂)	4250.0	686.20	8.95380	0.9×10^5

natural convection flow of an incompressible third-grade nanofluid are as follows [2, 3, 5, 14, 20].

The equation of motion is

$$\mu_{hmf} \frac{d^2 u}{dy^2} + 6\beta_3 \left(\frac{du}{dy} \right)^2 \frac{d^2 u}{dy^2} + (\beta_T \rho)_{hmf} (T - T_m) g = 0, \quad (2)$$

and the equation of energy is as follow:

$$K_{hmf} \frac{d^2 T}{dy^2} + \mu_{hmf} \left(\frac{dV}{dy} \right)^2 + 2\beta_3 \left(\frac{dV}{dy} \right)^4 = 0, \quad (3)$$

with the following boundary conditions:

$$\begin{aligned} u(y) &= 0, \\ \theta(y) &= T_1, \quad \text{at } y = -h, \\ u(y) &= 0, \\ \theta(y) &= T_2 \quad \text{at } y = h. \end{aligned} \quad (4)$$

The dimensionless variables [2]

$$\begin{aligned}
 u &= \frac{\bar{u}}{V_0}, \\
 y &= \frac{\bar{y}}{h}, \\
 \theta &= \frac{T - T_m}{T_1 - T_2},
 \end{aligned} \tag{5}$$

After removing the bar, we have

$$\frac{d^2 u}{dy^2} + \frac{6\beta}{A_2} \left(\frac{du}{dy} \right)^2 \frac{d^2 u}{dy^2} + \frac{A_3 A_1 \theta Gr}{A_2} = 0, \tag{6}$$

$$\frac{d^2 \theta}{dy^2} + \frac{PrEc}{A} (1 - \phi_1)^{2.5} (1 - \phi_2)^{2.5} \left(\frac{d\theta}{dy} \right)^2 + \frac{2\beta PrEc}{A} \left(\frac{d\theta}{dy} \right)^4 = 0, \tag{7}$$

And dynamic and thermal boundary conditions are

$$\begin{aligned}
 u(y) &= 0, \\
 \theta(y) &= -0.5, \quad \text{at } y = -1, \\
 u(y) &= 0, \\
 \theta(y) &= 0.5, \quad \text{at } y = 1.
 \end{aligned} \tag{8}$$

$$\begin{aligned}
 Pr &= \frac{\mu_f (\rho C_p)_f}{\rho_f k_f}, \\
 Ec &= \frac{V_0^2 \rho_f}{(T_1 - T_2) (\rho C_p)_f}, \\
 \beta &= \frac{6V_0^2 \beta_3}{h^2 \mu_f}, \\
 Gr &= \frac{(T_w - T_\infty) g (\rho \beta_T)_f}{h^2},
 \end{aligned} \tag{9}$$

where the dimensionless Grashof number (Gr), the Eckert number (Ec), Prandtl number (Pr), and the non-Newtonian viscosity (β).

$$A_1 = \frac{\rho_{hnf}}{\rho_f} = \left[(-\phi_2 + 1) \left\{ (1 - \phi_1) + \frac{\rho_{s_1}}{\rho_f} \phi_1 \right\} + \phi_2 \frac{\rho_{s_2}}{\rho_f} \right],$$

$$A_2 = \mu_{hnf} = \frac{\mu_f}{(1 - \phi_1)^{2.5} (1 - \phi_2)^{2.5}},$$

$$A_3 = \frac{(\beta_T)_{hnf}}{(\beta_T)_f} = \phi_2 \frac{(\beta_T)_{s_2}}{(\beta_T)_f} + \left[(1 - \phi_1) + \phi_1 \frac{(\beta_T)_{s_1}}{(\beta_T)_f} \right] (1 - \phi_2),$$

$$A = \frac{k_{hnf}}{k_{nf}} = \frac{2k_{nf} + 2\phi_1(k_{s_1} - k_{nf}) + k_{s_1}}{2k_{nf} - \phi_1(k_{s_1} - k_{nf}) + k_{s_1}},$$

$$\frac{k_{nf}}{k_f} = \frac{2k_f + 2\phi_2(k_{s_2} - k_f) + k_{s_2}}{2k_f - \phi_2(k_{s_2} - k_f) + k_{s_2}},$$

(10)

where ρ_{hnf} , k_{hnf} , μ_{hnf} , $(\beta_T)_{hnf}$, $(\rho C_p)_{hnf}$, ϕ_1 , and ϕ_2 denote the density, thermal conductivity, viscosity, thermal expansion coefficient, specific heat, Fe_3O_4 nanoparticles volume fraction, and TiO_2 nanoparticles volume fraction of hybrid nanofluids, respectively. [53].

4. Formulation of the Crisp Problem into a Fuzzy Problem Using FDEs

The velocity and temperature of nanoparticles are affected by small changes in their volume fraction. Some researchers take the nanoparticles volume fraction in this range [0.01–0.04], implying that fluid flow is solely dependent on these values. Then, due to the fixed crisp values of the volume fraction of nanoparticles, uncertainty develops.

Since ϕ_1 representing the volume fraction of Fe_3O_4 and ϕ_2 represents the volume fraction of TiO_2 , so, in a fuzzy environment, it is preferable to address a complex situation by accepting both volume fractions as FN.

For fuzzy solutions, equations (6)–(8) can be converted into FDE using α – cut approach. So, according to Definitions 4 and 5, we have

$$\begin{aligned} \frac{d^2}{dy^2} [u_1(y, \alpha), u_2(y, \alpha)] + \frac{6\beta}{A_2} \frac{d^2}{dy^2} [u_1(y, \alpha), u_2(y, \alpha)] \left(\frac{d}{dy} [u_1(y, \alpha), u_2(y, \alpha)] \right)^2 \\ + \frac{A_1 A_3 \text{Gr}}{A_2} [\theta_1(y, \alpha), \theta_2(y, \alpha)] = 0, \end{aligned} \quad (11)$$

$$\begin{aligned} \frac{d^2}{dy^2} [\theta_1(y, \alpha), \theta_2(y, \alpha)] + \frac{\text{PrEc}}{A} (1 - \phi_1)^{2.5} (1 - \phi_2)^{2.5} \left(\frac{d}{dy} [u_1(y, \alpha), u_2(y, \alpha)] \right)^2 \\ + \frac{2\beta \text{PrEc}}{A} \left(\frac{d}{dy} [u_1(y, \alpha), u_2(y, \alpha)] \right)^4 = 0, \end{aligned} \quad (12)$$

$$\begin{aligned} u(y, \alpha) &= 0, \\ \theta(y, \alpha) &= -0.5, \quad \text{at } y = -1, \\ u(y, \alpha) &= 0, \\ \theta(y, \alpha) &= 0.5, \quad \text{at } y = 1, \end{aligned} \quad (13)$$

where $0 \leq \alpha \leq 1$ $u_1(y, \alpha)$ is the lower bound and $u_2(y, \alpha)$ is the upper bound of fuzzy velocity profiles. Similarly, the fuzzy temperature profiles are $\theta(y, \alpha) = [\theta_1(y, \alpha), \theta_2(y, \alpha)]$, $0 \leq \alpha \leq 1$.

Table 2 presents the crisp values and TFNs of these FNs. The TFN defined the variation of FN at each α -cut. The TFNs are used to define the triangular MFs of the FNs which is ranging from 0 to 1, see Figure 1. This investigated range is commonly used to develop the aforementioned problem.

Now, we present a boundary value problem solver numerical procedure for controlling crisp differential equations (equations (6)–(8)) and FDEs (equations (11)–(13)) with boundary conditions, which are called `bvp4c`. It is a Lobatto IIIa formula with three stages based on the finite-difference algorithm. It has a collocation polynomial, and in $[a, b]$, the collocation formula yields a sixth-order accurate uniform C1 continuous solution. For error control and mesh selection, the continuous solutions residual is employed. The aforementioned ODEs are transformed into a first-order system as follows:

Let

$$\begin{aligned} u(y) &= m(1), \\ u'(y) &= m'(1) = m(2), \end{aligned} \quad (14)$$

$$\begin{aligned} u''(y) &= m'(2), \\ m'(2) &= \frac{-A_1 A_3 m(3) \text{Gr}}{1 + (1/A_2) 6\beta (m(2))^2}, \end{aligned} \quad (15)$$

$$\begin{aligned} \theta(y) &= m(3), \\ \theta'(y) &= m'(3) = m(4), \\ \theta''(y) &= m'(4), \end{aligned} \quad (16)$$

$$\begin{aligned} m'(4) &= \frac{1}{A} \text{PrEc} \left[(1 - \phi_1)^{-2.5} (1 - \phi_2)^{-2.5} (m(4))^2 \right. \\ &\quad \left. + 2\beta (m(4))^4 \right], \end{aligned} \quad (17)$$

and boundary conditions are

$$\begin{aligned} m_a(1) &= 0, \\ m_a(3) &= -0.5 \quad \text{at } y = -1, \\ m_b(3) &= 0, \\ m_b(3) &= 0.5, \quad \text{at } y = 1. \end{aligned} \quad (18)$$

For the required solution, equations (14) to (18) are coded in MATLAB software.

5. Results and Discussion

The SA is chosen as host fluid and $\text{Fe}_3\text{O}_4 + \text{TiO}_2$ are hybrid nanoparticles added into the base fluid to improve the rate of heat transfer between two vertical flat plates. The numerical solutions of governing coupled nonlinear DEs are obtained via the built-in MATLAB numerical technique `bvp4c`. The effect of thermo-physical parameters, such as Eckert number (Ec), Prandtl number (Pr), viscous dissipation parameter, Grashof number (Gr), third-grade fluid parameter (β), and nanoparticles volume fraction ϕ_1 and ϕ_2 on velocity and temperature fields are drawn in Figures 3–9.

Tables 3 and 4 show the comparison of velocity and temperature fields at $\phi_1 = \phi_2 = 0$, $\beta = 0.5$, $\text{Gr} = \text{Pr} = \text{Ec} = 1$, with studies by Ziabakhsh and Domairry [3], Manshoor et al. [5], and Biswal et al. [18, 19]. For the validation, the current study findings were found to be in excellent agreement.

Figure 3 displays the influence of the Prandtl number (Pr) on the velocity and temperature fields while other physical parameters are fixed. The velocity and temperature fields of the hybrid nanofluid rise as Pr increases due to upsurges in the thickness of the boundary layer.

The impact of the viscous dissipation parameter (Ec) on the velocity and temperature fields is demonstrated in Figure 4. It can be observed that the velocity and temperature of the hybrid nanofluid enhance with growing the values of Ec. When Ec increases, the dissipation of heat on the boundary layer region increases and also the heat transfer rate increases. The impact of third-grade fluid parameter (β) on the velocity and temperature field

TABLE 2: TFNs of fuzzy nanoparticles of volume fraction.

Fuzzy number	Crisp value	TFN	α - cut approach
ϕ_1	[0.01-0.04]	[0, 0.05, 0.1]	$[0.05\alpha, 0.1 - 0.05\alpha], \alpha \in [0, 1]$
ϕ_2	[0.01-0.04]	[0, 0.05, 0.1]	$[0.05\alpha, 0.1 - 0.05\alpha], \alpha \in [0, 1]$

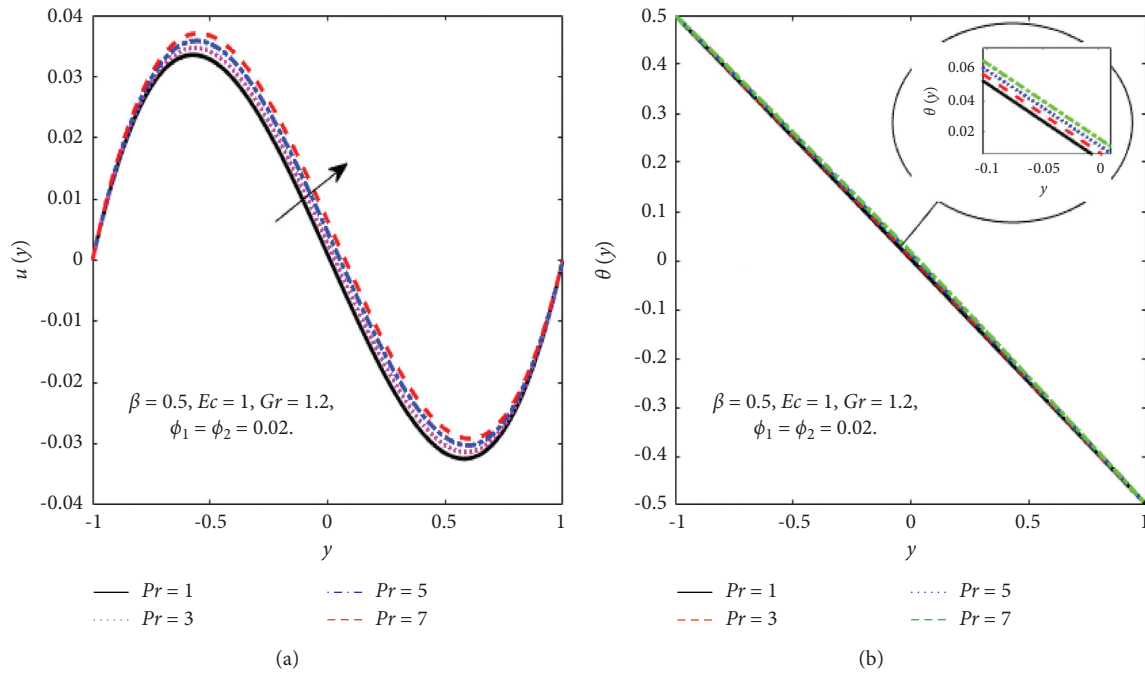


FIGURE 3: Effect of Pr on the $u(y)$ (a) and $\theta(y)$ (b).

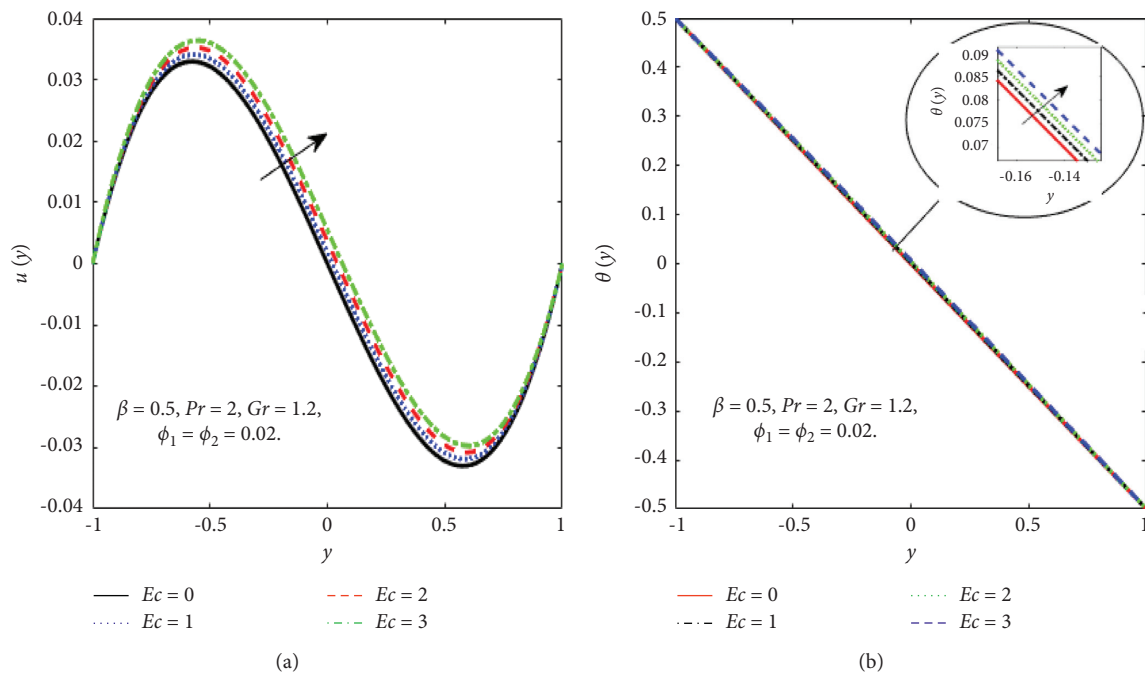


FIGURE 4: Effect of Ec on the $u(y)$ (a) and $\theta(y)$ (b).

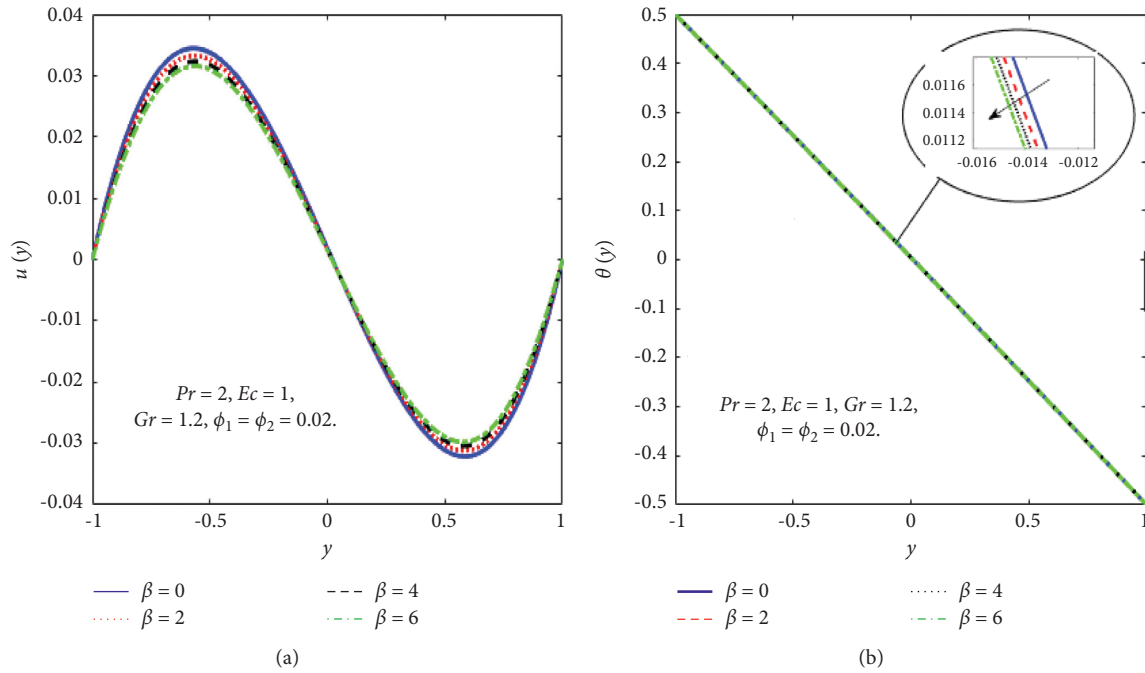


FIGURE 5: Effect of β on the $u(y)$ (a) and $\theta(y)$ (b).

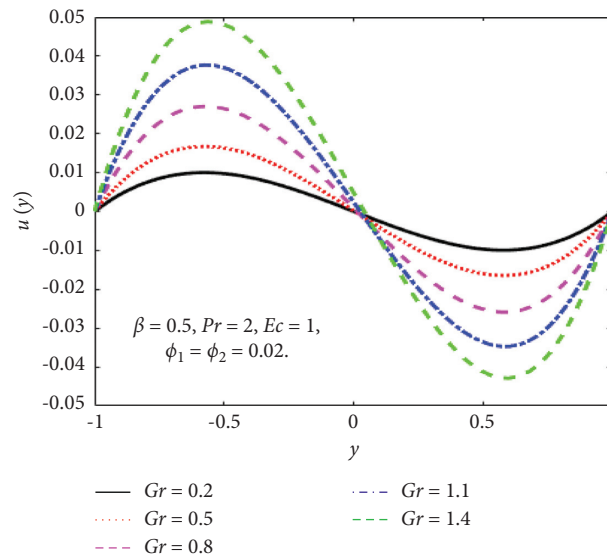


FIGURE 6: Effect of Gr on the $u(y)$.

of the hybrid nanofluid is examined in Figure 5. In Figure 5(a), when β increases, then the velocity declines in the region $-0.9 < y < -0.1$ and it is increasing in the region $0.4 < y < 0.8$. The reason for this is that when the viscosity of the hybrid nanofluid increases, the boundary layer thickens and the velocity declines. In Figure 5(b), when the value of β increases, the temperature falls. The temperature profile shows very small variations on large values of β because the rate of shear increases and decreases in the boundary layer thickness. The impact of

buoyancy forces (Gr) on the velocity profile is portrayed in Figure 6. It can be seen that when Gr is amplified, then the velocity profile displays an increasing trend. Physically, large values of Gr boost the buoyancy force, resulting in a higher thermal force through the use of the viscous force and hence there is an upsurge in hybrid nanofluid velocity. Figures 7 and 8 demonstrate the impact of hybrid nanoparticles volume fraction (ϕ_1, ϕ_2) on velocity and temperature fields. These profiles are plotted for hybrid nanofluids ($(Fe_3O_4 + TiO_2)/SA$). In

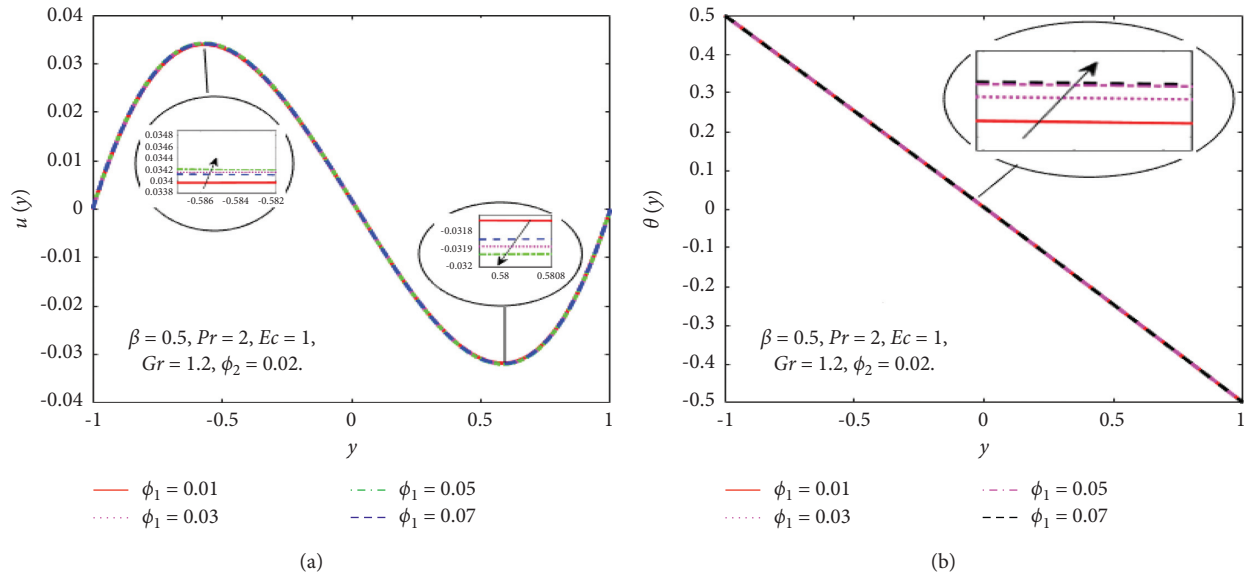


FIGURE 7: Effect of ϕ_1 on the $u(y)$ (a) and $\theta(y)$ (b).

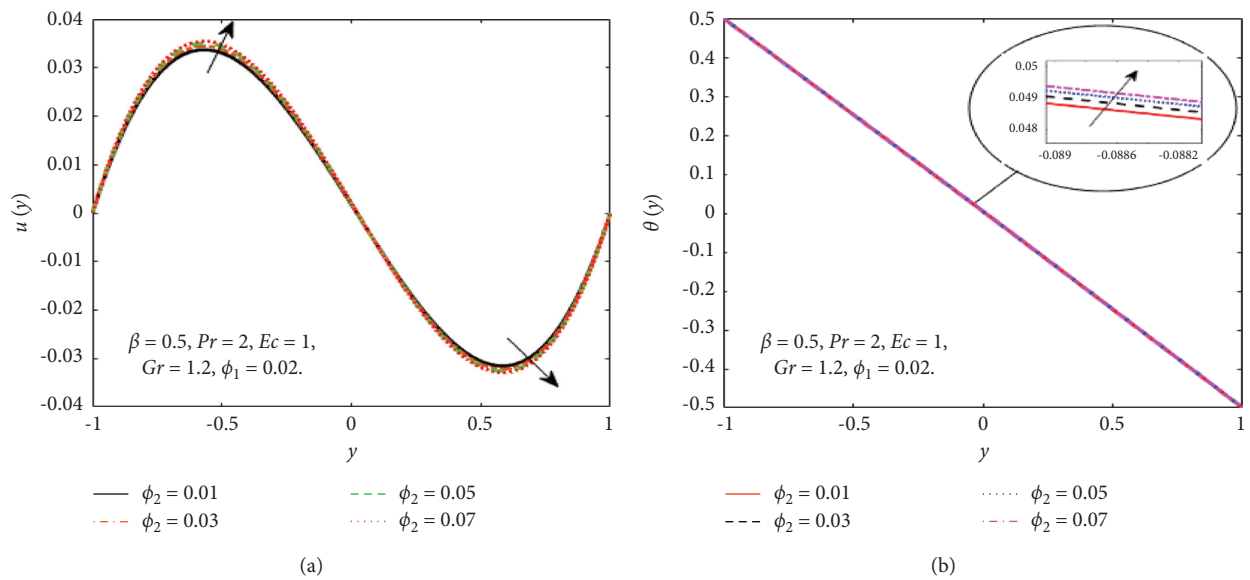


FIGURE 8: Effect of ϕ_2 on the $u(y)$ (a) and $\theta(y)$ (b).

Figure 7(a), the velocity of hybrid nanofluid increases with an increase in ϕ_1 and Figure 7(b) shows the temperature profile increases because of increased heat transfer at $\phi_2 = 0.01$. The reason is that the friction of solid particles decreases the host fluid viscosity. Similarly, Figure 8(a) shows that the velocity of hybrid nanofluids increases with an increase in ϕ_2 and Figure 8(b) shows that the temperature profile increases because of increased heat transfer at $\phi_1 = 0.01$. Physically, the intermolecular forces between the particles of hybrid nanofluids become weaker, and consequently, the hybrid nanofluid velocity accelerates. Further, it is detected that the thermal boundary layer thickness increases because

the temperature profile increases due to higher values of hybrid nanoparticle's volume fraction. Figure 9 represents the comparison of nanofluids Fe_3O_4/SA and TiO_2/SA and hybrid nanofluid $((Fe_3O_4 + TiO_2)/SA)$ at $\phi_1 = \phi_2 = 0.04$. The velocity and temperature profiles of Fe_3O_4/SA are calculated at $\phi_1 = 0.04$ and $\phi_2 = 0$, whereas the velocity and temperature profiles of TiO_2/SA were calculated at $\phi_2 = 0.04$ and $\phi_1 = 0$. The velocity and temperature profiles of Fe_3O_4 are greater than the velocity and temperature profiles of TiO_2 . Also, the velocity and temperature profiles of hybrid nanofluids $((Fe_3O_4 + TiO_2)/SA)$ are greater than Fe_3O_4 and TiO_2 . Physically, this is correct because Fe_3O_4 has a higher heat

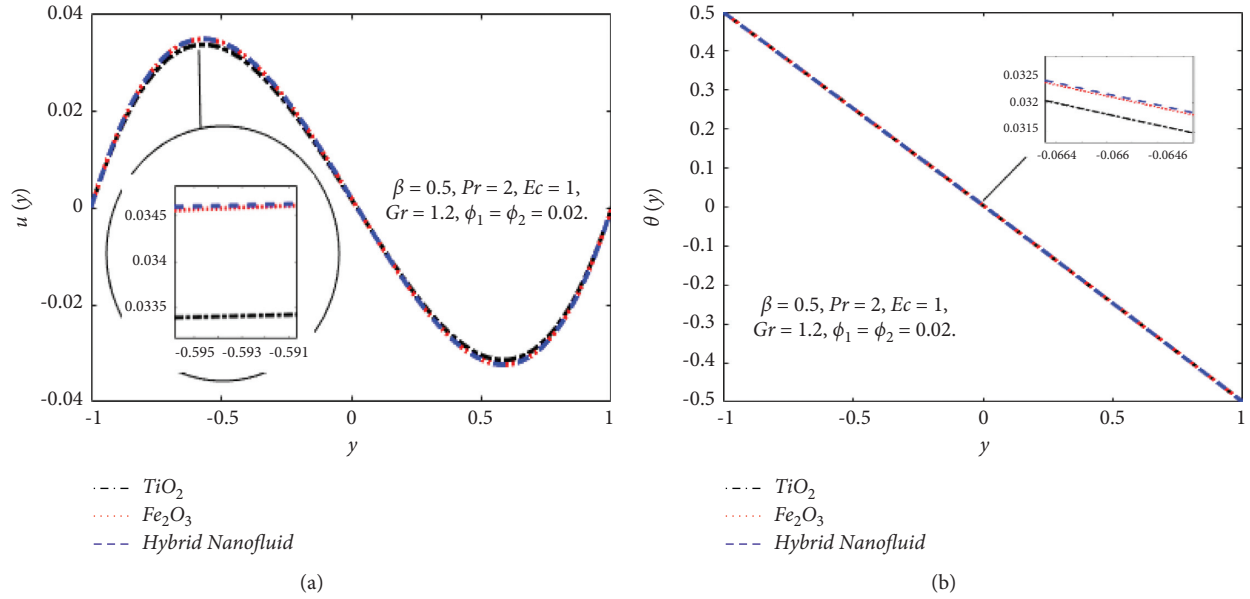


FIGURE 9: Comparison of Fe_3O_4 , TiO_2 , and hybrid nanofluids on the $u(y)$ (a) and $\theta(y)$ (b).

TABLE 3: Comparison of velocity profile when $\phi_1 = \phi_2 = 0, \beta = 0.5, Gr = Pr = Ec = 1$, with the existing result for regular fluid.

Y	Present results (bvp4c)	Ziabakhsh and Domairry [3] (HAM)	Biswal et al. [18] (HPM)	Biswal et al. [19] (GM)	Manshoor et al. [5] (VPM)
-1	0	0	0	0	0
-0.8	0.02244430	0.02391937	0.02416171	0.02368610	0.033923604
-0.6	0.03122643	0.03217274	0.03262933	0.03170120	0.032183540
-0.4	0.02712964	0.02840695	0.02901579	0.02794809	0.027143138
-0.2	0.01603016	0.01661778	0.01731154	0.01632954	0.016274634
0	-0.00002592	0.00080780	0.00152009	0.00074834	0.000922405
0.2	-0.01448865	-0.01508225	-0.01441546	-0.01489272	-0.015143973
0.4	-0.02683483	-0.02710348	-0.026554082	-0.02669087	-0.028257013
0.6	-0.03070475	-0.03122988	-0.03081889	-0.03074332	-0.031223835
0.8	-0.02313980	-0.02342875	-0.02320304	-0.02314729	-0.023274354
1	0	0	0	0	0

TABLE 4: Comparison of temperature profile when $\phi_1 = \phi_2 = 0, \beta = 0.5, Gr = Pr = Ec = 1$, with the existing result for regular fluid.

Y	Present results (bvp4c)	Ziabakhsh and Domairry [3] (HAM)	Biswal et al. [18] (HPM)	Biswal et al. [19] (GM)	Manshoor et al. [5] (VPM)
-1	0.49999999	0.49999999	0.49987599	0.50000000	0.500000000
-0.8	0.40009178	0.40073588	0.40157624	0.40097357	0.400246306
-0.6	0.30116719	0.30117737	0.30269966	0.30172607	0.309367078
-0.4	0.20863343	0.20159090	0.20321740	0.20225927	0.201548465
-0.2	0.10108217	0.10192749	0.10350286	0.10257493	0.101925345
0	-0.00299024	0.00206049	0.00350177	0.00267484	0.002174325
0.2	-0.09501960	-0.09807006	-0.09677903	-0.09743924	-0.098174536
0.4	-0.19536608	-0.19840851	-0.19733286	-0.19776553	-0.198546725
0.6	-0.29679969	-0.29882852	-0.29812783	-0.29830227	-0.298765434
0.8	-0.39825954	-0.39927474	-0.39909138	-0.39904768	-0.400465233
1	-0.49999999	-0.500000000	-0.50012401	-0.50000000	-0.500000000

conductivity than TiO_2 . However, the temperature profile shows the same behaviour as that of the velocity profile, as Fe_3O_4 has larger thermal conductivity than TiO_2 . As a result, Fe_3O_4 conducts more heat than TiO_2 and is less

dense, resulting in Fe_3O_4 having a higher temperature than TiO_2 . Considering these factors, this study recommends using Fe_3O_4 to improve heat transmission since Fe_3O_4 conducts more heat and is more stable than TiO_2 .

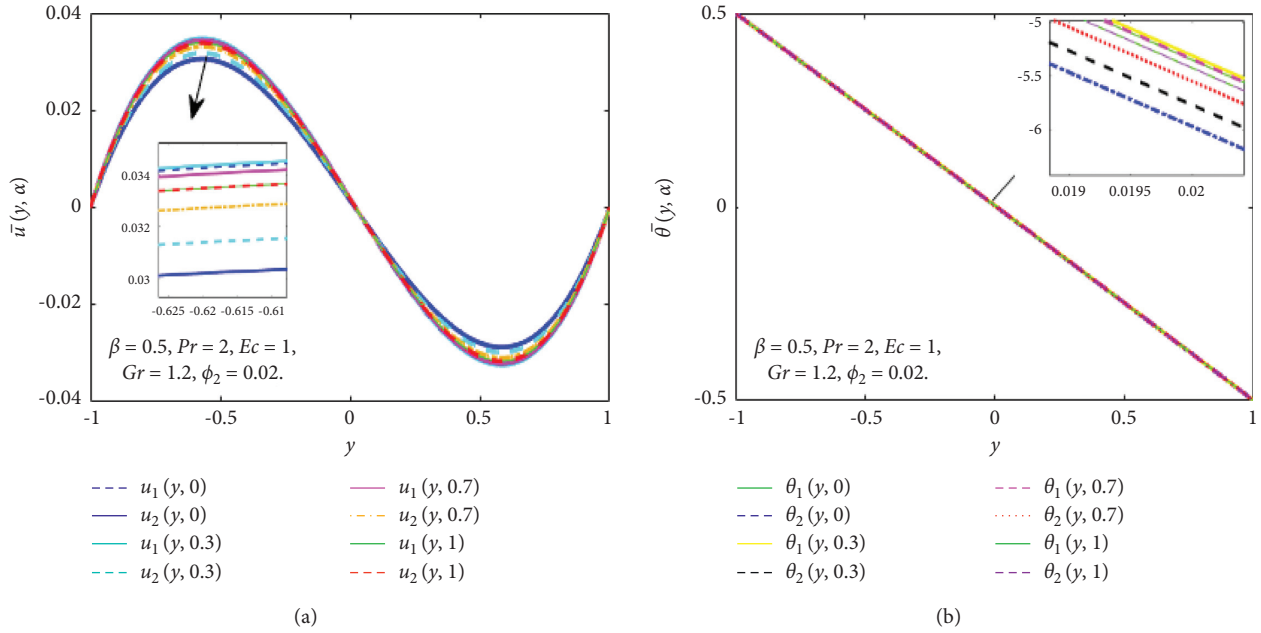


FIGURE 10: Effect of ϕ_1 on the fuzzy velocity (a) and temperature (b) profiles if ϕ_1 is a TFN.

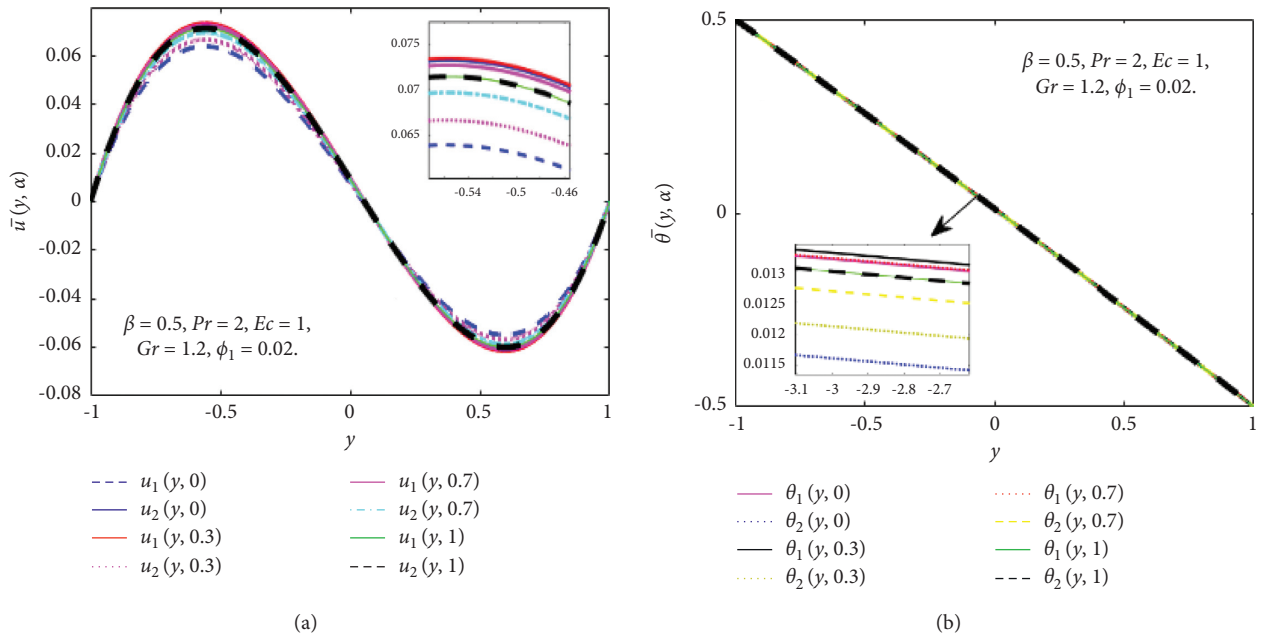


FIGURE 11: Effect of ϕ_2 on the fuzzy velocity (a) and temperature (b) profiles if ϕ_2 is a TFN.

Now, we discuss the nanoparticles volume fraction of Fe_3O_4 (ϕ_1) and TiO_2 (ϕ_2) in a fuzzy environment. The nanoparticles volume fraction ϕ_1 and ϕ_2 are said to be TFN, as shown in Table 1, and analyzed by α -cut approach ($0 \leq \alpha \leq 1$), as discussed in Section 3.4 in detail.

Figures 10 and 11 show the nanoparticles volume fraction of Fe_3O_4 (ϕ_1) and TiO_2 (ϕ_2) considered as TFNs (see in Table 2) and then the $\bar{u}(y, \alpha)$ and $\bar{\theta}(y, \alpha)$ are controlled by α -cut for some particular values of α -cut ($\alpha = 0, 0.3, 0.7, 1$). In Figure 10, when ϕ_1 is a TFNs, then the

$\bar{u}(y, \alpha)$ and $\bar{\theta}(y, \alpha)$ convert into lower and upper bounds of the velocity and temperature fields. When α -cut = 0, $u_1(y, \alpha)$ and $\theta_1(y, \alpha)$ represent the nanofluid while $u_2(y, \alpha)$ and $\theta_2(y, \alpha)$ represent hybrid nanofluid at $\phi_2 = 0.04$. When α -cut increases the width between $u_1(y, \alpha)$ and $u_2(y, \alpha)$ decreases and at α -cut = 1, they coherent with one another. It is noted that the width between $u_1(y, \alpha)$ and $u_2(y, \alpha)$ is very less, so the vagueness is less. Similarly, in the case of the $\bar{\theta}(y, \alpha)$, as α rises, the width between $\theta_1(y, \alpha)$ and $\theta_2(y, \alpha)$ reduces, and $\alpha = 1$, they are coherent with one another. It is

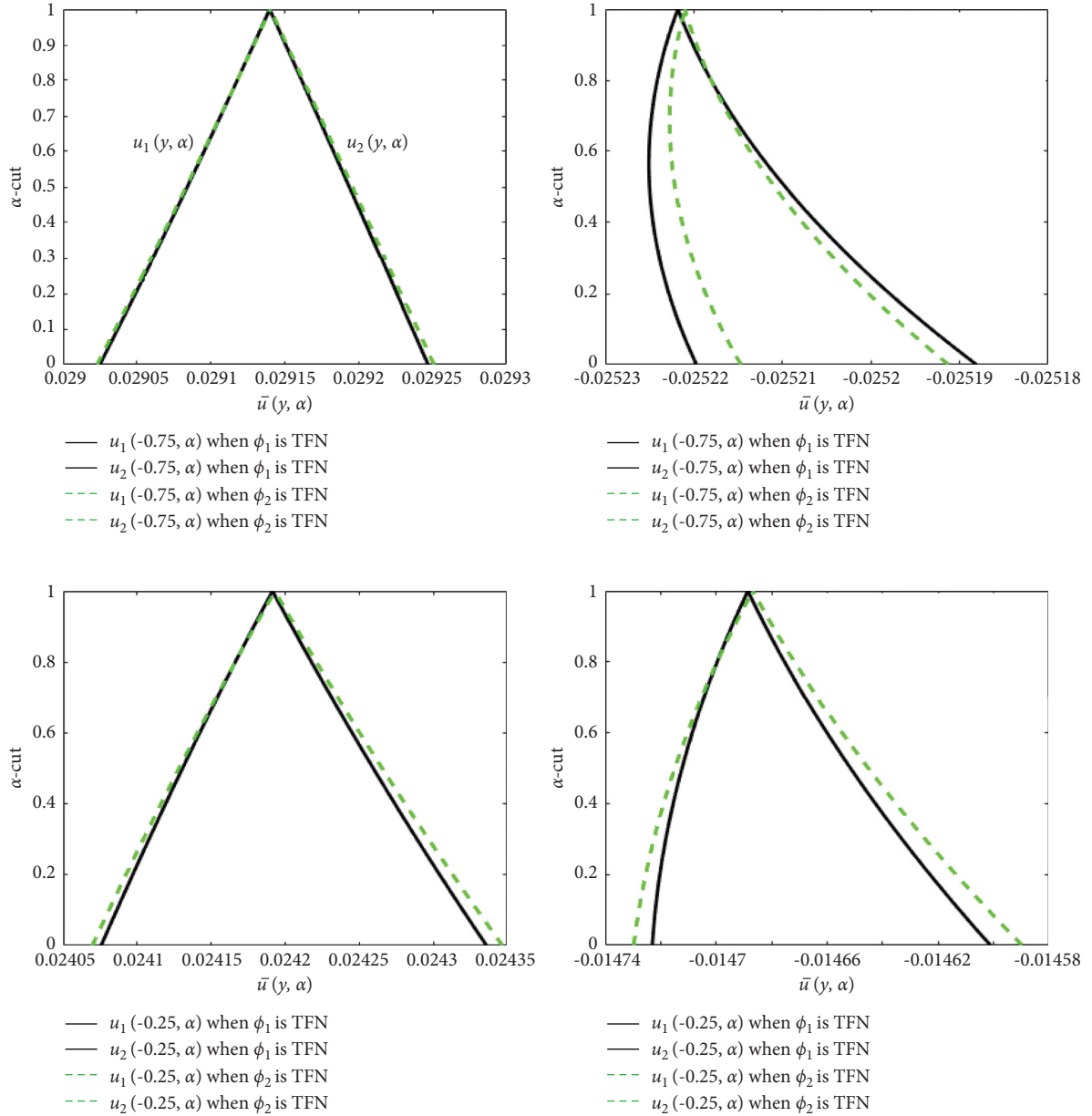


FIGURE 12: Effects of TFNs ϕ_2 and ϕ_2 on the fuzzy velocity profile. (a) Fuzzy velocity at $y = -0.75$. (b) Fuzzy velocity at $y = 0.75$. (c) Fuzzy velocity at $y = -0.25$. (d) Fuzzy velocity at $y = 0.25$.

vital to keep in mind that the width between $\theta_1(y, \alpha)$ and $\theta_2(y, \alpha)$ is quite narrow, indicating that the uncertainty is very less. Consequently, in Figure 11, when ϕ_2 is a TFN, then the $\bar{u}(y, \alpha)$ and $\bar{\theta}(y, \alpha)$ convert into $u_1(y, \alpha)$, $u_2(y, \alpha)$, $\theta_1(y, \alpha)$, and $\theta_2(y, \alpha)$. It is essential to note that the width between the lower and upper bounds of velocity and temperature fields is very narrow, which indicates that the uncertainty is minimal.

Figures 12 and 13 show the triangular membership functions of the $\bar{u}(y, \alpha)$ and $\bar{\theta}(y, \alpha)$ for various values of y .

In these diagrams, we investigated two different cases. The black lines represent the case where ϕ_1 is used as the TFN and $\phi_2 = 0.04$. The green and red dashed lines indicate the representation of ϕ_2 as TFN, whereas $\phi_1 = 0.04$. The horizontal axis displays the $\bar{u}(y, \alpha)$ and $\bar{\theta}(y, \alpha)$ for varying y , while the vertical axis displays the membership values of the $\bar{u}(y, \alpha)$ and $\bar{\theta}(y, \alpha)$ for varying α -cut. From Figure 12, it can be seen that the width between $u_1(y, \alpha)$ and $u_2(y, \alpha)$ is less, therefore the uncertainty is less for numerous values of y . The width between $\theta_1(y, \alpha)$ and $\theta_2(y, \alpha)$ is moderately

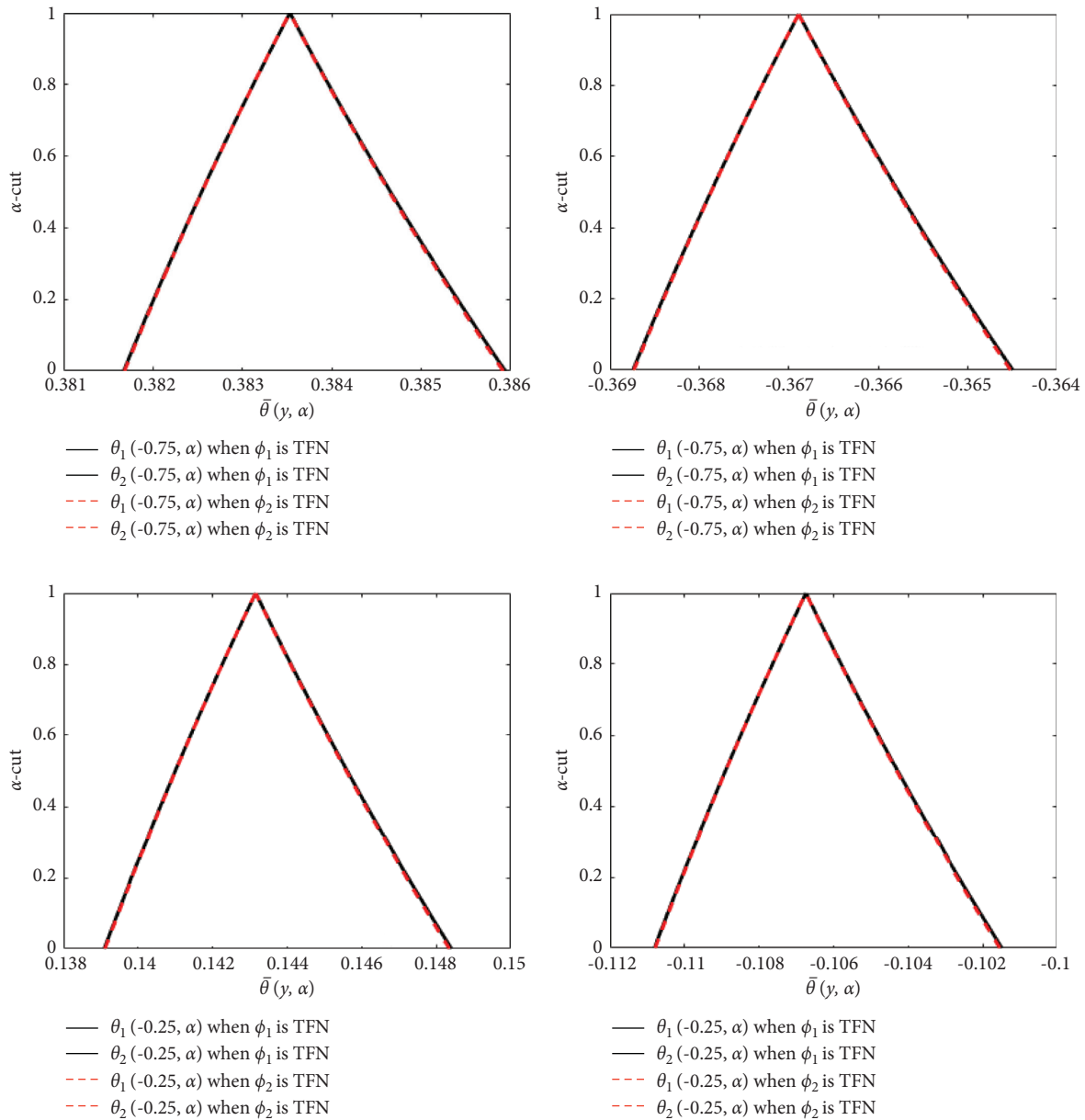


FIGURE 13: Effects of TFNs ϕ_1 and ϕ_2 on the fuzzy temperature profile. (a) Fuzzy temperature at $y = -0.75$. (b) Fuzzy temperature at $y = 0.75$. (c) Fuzzy temperature at $y = -0.25$. (d) Fuzzy temperature at $y = 0.25$.

slight in Figure 13, demonstrating that the impreciseness is neglectable for various values of y . As a result, the uncertain parameters are controlled through TFNs.

6. Conclusion

The current study focused on the natural convection flow of third-grade $(Fe_3O_4 + TiO_2)/SA$ hybrid nanofluid across vertical parallel plates in a fuzzy environment. The impacts of the Eckert number (Ec), the non-Newtonian viscosity (β), Prandtl number (Pr), Grashof number (Gr), and nanoparticles volume fraction (ϕ_1, ϕ_2) on the temperature and velocity profiles have been studied for $(Fe_3O_4 + TiO_2)/SA$

hybrid nanofluid. The volume fractions of nanoparticles of Fe_3O_4 (ϕ_1) and TiO_2 (ϕ_2) are considered as TFNs with the help of α -cut ($0 \leq \alpha \leq 1$) which control fuzziness. For various values of y , triangular membership plots of fuzzy velocity and temperature profiles were also examined. The following significant finding comes from this investigation:

The velocity and temperature profiles rise as the values of Pr , Gr , and Ec increase, whereas the velocity and temperature profiles decrease when the value β increases.

The rate of heat transfer upsurges by growing volume fractions of nanoparticles ϕ_1 and ϕ_2 .

The present results obtained from numerical technique via $bvp4c$ are found to be in excellent agreement as compared to existing results.

The hybrid nanofluid ($Fe_3O_4 + TiO_2$)/SA shows a higher heat transfer rate as compared to nanofluids Fe_3O_4 /SA and TiO_2 /SA.

The results indicate that the crisp solution is always in-between the upper and lower solutions when α – cut to increase from 0 to 1.

The sensitivity of the assumed TFN is held influenced by the unfluctuating width of the fuzzy velocity or temperature.

According to the triangular membership plots, the uncertain width of the fuzzy velocity and temperature is less, so the assumed TFNs are less sensitive. Finally, the TFN is represented visually for better understanding. As a result, the TFNs may be used to different heat transfer problems.

Data Availability

No data were used in this article.

Conflicts of Interest

The authors declare that they have no conflicts of interest.

Acknowledgments

The authors extend their appreciation to the Deanship of Scientific Research at Majmaah University for funding this work under project no. R.G.P. 2019–4.

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